

Integrable Systems: Lecture 4

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Contents

In this lecture, we continued the study of the scattering theory for the Schrödinger operator. We followed the exposition from Section 1.2: *Elements of scattering theory for the Schrödinger operator* from [Dub2009].

References

[Dub2009] B. A. Dubrovin, *Integrable systems and Riemann surfaces*, 2009. lecture notes, available at the author's web-page.

$$a_1' = \frac{1}{2ik} e^{-ikx} f \quad a_2' = -\frac{1}{2ik} e^{ikx} f$$

$$\downarrow$$

$$a_1(x) = a_1^0 + \frac{1}{2ik} \int_{x_0}^x f(y) e^{-iky} dy$$

$$a_2(x) = a_2^0 - \frac{1}{2ik} \int_{x_0}^x f(y) e^{iky} dy$$

$$\downarrow$$

$$\Psi(x) = \frac{1}{2ik} \int_{x_0}^x e^{ik(x-y)} u(y) \Psi(y) dy - \frac{1}{2ik} \int_{x_0}^x e^{ik(y-x)} u(y) \Psi(y) dy + a_1^0 e^{ikx} + a_2^0 e^{-ikx}$$

We want: $\Psi \sim e^{ikx}$ as $x \rightarrow +\infty$ | We set $x_0 = +\infty$
 $a_1^0 = 1$ $a_2^0 = 0$

$$\Psi(x) = e^{ikx} - \int_x^{+\infty} \frac{\sin k(x-y)}{k} u(y) \Psi(y) dy$$

$$\Psi = f_0 + f_1 + f_2 + \dots$$

$$f_0 = e^{ikx}$$

$$f_{n+1}(x) = - \int_x^{+\infty} \frac{\sin k(x-y)}{k} u(y) f_n(y) dy$$

We need to prove that the sum is uniformly convergent.

$$|f_n(x)| \leq \frac{1}{n!} U^n(x), \quad U(x) = \frac{1}{k} \int_x^{+\infty} |u(y)| dy$$

Proved by induction.

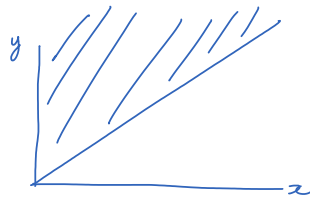
$$n=0: |f_0(x)| = |e^{ikx}| = 1 \checkmark$$

$$n=1: f_1(x) = - \int_x^{+\infty} \frac{\sin k(x-y)}{k} u(y) f_0(y) dy$$

$$|f_1(x)| \leq \int_x^{+\infty} \frac{|\sin k(x-y)|}{|k|} |u(y)| |f_0(y)| dy$$

$$\leq \frac{1}{|k|} \int_x^{+\infty} |u(y)| dy \checkmark$$

$$n=2: |f_2(x)| \leq \int_x^{+\infty} \frac{1}{|k|} |u(y)| |f_1(y)| dy \leq \int_x^{+\infty} |u(y)| U(y) dy$$



$$\begin{aligned}
 |\Psi(z) - e^{iz}| &\leq |f_1| + |f_2| + |f_3| + \dots \\
 &= \sum_{n=1}^{\infty} \frac{1}{n!} O^n(x) \\
 &= |e^{O(z)} - 1| \xrightarrow{\text{as } z \rightarrow +\infty} 0
 \end{aligned}$$

Thus, we proved that Ψ exists and is unique
 and $\Psi \sim e^{iz}$ as $z \rightarrow +\infty$
 $\Psi_2 = \overline{\Psi_1}$ - existence and uniqueness hold for Ψ_1 as well

Now we prove that Ψ_2 is analytic for $\text{Im} k > 0$.

Denote: $\Psi(z) = e^{iz} \chi(z)$

$$\chi(z) = 1 - \int_z^{+\infty} \frac{1 - e^{2ib(y-z)}}{2ib} u(y) \chi(y) dy$$

$\left| \frac{2ib(y-z)}{e^{2ib(y-z)}} \right| = e^{-2\text{Im}(b)(y-z)} \rightarrow 0$
 for $|k| \rightarrow \infty$

$$\Psi_2 e^{-ibz} = \chi = 1 + O\left(\frac{1}{z}\right) \text{ and it is analytic in } \mathbb{C}$$

□

We proved that for each k ,

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there is a unique solution of
$$L\psi = \lambda^2 \psi$$

with $\psi \sim e^{ikx}$ as $x \rightarrow +\infty$

Similarly, we can prove existence and uniqueness of φ_1, φ_2 :

$\varphi_1 \sim e^{-ikx}$
 $\varphi_2 \sim e^{ikx}$ as $x \rightarrow -\infty$

We can prove that φ_1 extends analytically to $\text{Im} k > 0$

Denote: $\varphi := \varphi_2$ we have $\varphi_1 = \bar{\varphi}$
 $\varphi := \varphi_1$ we have $\varphi_2 = \bar{\varphi}$

Lemma $(\psi, \bar{\varphi})$ and $(\varphi, \bar{\psi})$ are two bases in the space of solutions. ($k \neq 0$)

Proof: $W(\psi, \bar{\varphi})$ does not depend on x

$$\lim_{x \rightarrow +\infty} W(\psi, \bar{\varphi}) = 2ik \neq 0 \quad \square$$

$$(\bar{\Psi}, \Psi) = (\Psi_1, \Psi_2)$$

$$(\Psi, \bar{\Psi}) = -(\Psi_1, \Psi_2)$$

$$\Psi(x, z) = a(z) \bar{\Psi}(x, z) + b(z) \Psi(x, z)$$

$$\bar{\Psi}(x, z) = \bar{a}(z) \Psi(x, z) + \bar{b}(z) \bar{\Psi}(x, z)$$

the transition matrix is $A = \begin{pmatrix} a(z) & b(z) \\ \bar{b}(z) & \bar{a}(z) \end{pmatrix}$

Lemma $|a(z)|^2 - |b(z)|^2 = 1$

Proof

$$\det A = 1 \quad \text{in } \mathbb{C}$$

$$\begin{pmatrix} \Psi \\ \bar{\Psi} \end{pmatrix} = A \begin{pmatrix} \bar{\Psi} \\ \Psi \end{pmatrix}$$

□

Lemma $a(z)$ can be analytically extended to $\Im z > 0$.

Proof.

$$W(\Psi, \Psi) = W(\Psi, a\bar{\Psi} + b\Psi) =$$

$$= aW(\Psi, \bar{\Psi}) + bW(\Psi, \Psi) \rightarrow 0$$

$$= a(z)W(\Psi, \bar{\Psi}) = -2ik a(z)$$

$$\Gamma W(f, g) = \det \begin{pmatrix} f & g \\ f' & g' \end{pmatrix}$$

$$\downarrow$$

$$a(z) = -\frac{i}{2k} W(\Psi, \Psi)$$

□