

# Integrable Systems: Lecture 7

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## Contents

In this lecture, we showed that the inverse scattering can be reduced to a linear problem. We followed the exposition from [Dub2009, Section 1.3: *Inverse scattering*] and [NMPZ1984, Section 1.1: *Elements of the scattering theory. The inverse problem in the quantum theory of scattering*].

## References

- [Dub2009] B. A. Dubrovin, *Integrable systems and Riemann surfaces*, 2009. lecture notes, available at the author's web-page.
- [NMPZ1984] S. Novikov, S. V. Manakov, L. P. Pitaevskii, and V. E. Zakharov, *Theory of solitons*, Contemporary Soviet Mathematics, Consultants Bureau [Plenum], New York, 1984. The inverse scattering method; Translated from the Russian.

# Inverse scattering

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$$3^\circ \quad u(x) = -2 \frac{d}{dx} K(x, x)$$

$$1^\circ \quad F(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} r(k) e^{ikx} dk + \sum_{s=1}^n \frac{b_s e^{-\alpha_s x}}{i a'(i\alpha_s)}$$

2<sup>o</sup> Solve the Gel'fand-Levitan-Marchenko equation  
in  $K = K(x, y)$ :

$$K(x, y) + F(x+y) + \int_{-\infty}^{\infty} K(x, z) F(z+y) dz = 0$$

$$\Psi(x, k) = a(k)\bar{\Psi}(x, k) + b(k)\Psi(x, k)$$

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$$\frac{e^{ikx}}{a(k)}$$

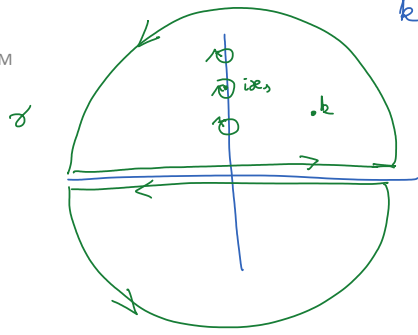
$$\underbrace{\Psi(x, k) \frac{e^{ikx}}{a(k)}}_{\substack{\text{analytic in} \\ \text{Im} k > 0, \\ \text{except having simple poles at } i\alpha_s}} = \underbrace{\bar{\Psi}(x, k) e^{ikx}}_{\text{analytic in } \text{Im} k < 0} + b(k)\Psi(x, k) e^{ikx}$$

$k \in \mathbb{R}$   
+  
+  
+  
+  
+  
+  
+  
+  
+  
+

Define:

$$\Phi(x, k) := \begin{cases} \frac{\Psi(x, k) e^{ikx}}{a(k)}, & \text{Im} k > 0 \\ \bar{\Psi}(x, k) e^{ikx}, & \text{Im} k < 0 \end{cases}$$

- properties of  $\Phi$ :
- simple poles at  $i\alpha_s$
  - "step" at  $k \in \mathbb{R}$
  - $\Phi \rightarrow 1$  as  $|k| \rightarrow \infty$
  - analytic everywhere else



$$\begin{aligned} \phi(x, b) - 1 &= \frac{1}{2\pi i} \int_{\gamma} \frac{\phi(x, k') - 1}{k' - b} dk' \\ &= \sum_{s=1}^n \frac{\Gamma_s(x)}{b - i\alpha_s} + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\gamma(k') \Psi(x, k') e^{ik'x}}{k' - b} dk' \end{aligned}$$

$\Gamma_s(x)$  - the residue of  $\phi(x, k)$  at  $i\alpha_s$

$$\begin{aligned} \Gamma_s(x) &= \frac{\Psi(x, i\alpha_s) e^{-\alpha_s x}}{a'(i\alpha_s)} = \frac{b_s \Psi(x, i\alpha_s) e^{-\alpha_s x}}{a'(i\alpha_s)} = \\ &= \frac{b_s \Psi(x, -i\alpha_s) e^{-\alpha_s x}}{a'(i\alpha_s)} = \frac{b_s \phi(x, -i\alpha_s) e^{-2\alpha_s x}}{a'(i\alpha_s)} \end{aligned}$$

Take the limit  $\gamma_m b \rightarrow 0_-$  in (4):

$$\Psi(x, -b) e^{ibx} = 1 + \sum \frac{\Gamma_s(x)}{b - i\alpha_s} + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\gamma(k') \Psi(x, k') e^{ik'x}}{k' - b} dk'$$

We have got  $(n+1)$  linear relations in  $(n+1)$  functions:  $\Gamma_1, \dots, \Gamma_n, \Psi$ .

we will plug  $\phi$  from (4) into this expression

$$\Gamma_s(x) = \frac{b_s e^{-2\alpha_s x}}{a'(i\alpha_s)} \left( 1 + i \sum_m \frac{\Gamma_m(x)}{\alpha_m + \alpha_s} + \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\gamma(k') \Psi(x, k') e^{ik'x}}{k' + i\alpha_s} dk' \right)$$

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We want to deduce the  
GLM equation.

$\bar{\Psi}(x, \xi) e^{i\xi x}$  - analytic for  $\text{Im} \xi < 0$

and  $\sim 1 + O(\frac{1}{\xi})$  as

$$\Rightarrow \bar{\Psi}(x, \xi) e^{i\xi x} - 1 = \int_0^{\infty} A(x, y) e^{-i\xi y} dy \quad |\xi| \rightarrow \infty$$

$$\begin{aligned} \bar{\Psi}(x, \xi) &= e^{-i\xi x} + \int_0^{\infty} A(x, y) e^{-i\xi(x+y)} dy \\ &= e^{-i\xi x} + \int_0^{\infty} A(x, y'-x) e^{-i\xi y'} dy' \quad \text{change: } y' = x+y \\ &= e^{-i\xi x} + \int_x^{\infty} A(x, y-x) e^{-i\xi y} dy \\ & \quad K = A(x, y-x) \\ \bar{\Psi}(x, \xi) &= e^{-i\xi x} + \int_x^{\infty} K(x, y) e^{-i\xi y} dy \end{aligned}$$