

MATH2065: INTRO TO PDES

Summer School 2012

Tutorial Solutions 5

1. Taking partial derivatives of u :

$$\begin{aligned}\frac{\partial u}{\partial t} &= -4\pi^2 k \sin 2\pi x e^{-4\pi^2 kt}, \\ \frac{\partial u}{\partial x} &= 2\pi \cos 2\pi x e^{-4\pi^2 kt}, \\ \frac{\partial^2 u}{\partial x^2} &= -4\pi^2 \sin 2\pi x e^{-4\pi^2 kt}.\end{aligned}$$

The last of these when multiplied by k gives $\frac{\partial u}{\partial t}$ as required.

2. Let $u(x, t) = \sin(x - ct)$. Taking partial derivatives, and utilising the chain rule,

$$\begin{aligned}\frac{\partial u}{\partial t} &= \cos(x - ct)(-c) = -c \cos(x - ct), \\ \frac{\partial^2 u}{\partial t^2} &= (-c)(-\sin(x - ct))(-c) = -c^2 \sin(x - ct), \\ \frac{\partial u}{\partial x} &= \cos(x - ct), \\ \frac{\partial^2 u}{\partial x^2} &= -\sin(x - ct).\end{aligned}$$

Therefore, multiplying u_{xx} by c^2 gives us u_{tt} , which means that

$$u_{tt} = c^2 u_{xx}.$$

Now consider the purported D'Alembert's solution

$$u(x, t) = f(x - ct) + g(x + ct).$$

Note that each of f and g are functions of *one* variable; what appears above is that the combination $x - ct$ has been fed into f , and the combination $x + ct$ into g . Now, using the chain rule,

$$\begin{aligned}\frac{\partial u}{\partial x} &= f'(x - ct) + g'(x + ct), \\ \frac{\partial^2 u}{\partial x^2} &= f''(x - ct) + g''(x + ct), \\ \frac{\partial u}{\partial t} &= f'(x - ct)(-c) + g'(x + ct)(+c) = c(-f'(x - ct) + g'(x + ct)), \\ \frac{\partial^2 u}{\partial t^2} &= c(-f''(x - ct)(-c) + g''(x + ct)(+c)) = c^2(f''(x - ct) + g''(x + ct)).\end{aligned}$$

and hence $u_{tt} = c^2 u_{xx}$. When applying the chain rule above, what we have done in each instance is given by the following example:

$$\begin{aligned}\frac{\partial}{\partial t} f(x - ct) &= (\text{Derivative of } f \text{ with respect to whatever's inside}) \\ &\quad \times (\text{Derivative of what's inside with respect to } t) \\ &= f'(x - ct) \frac{\partial}{\partial t} (x - ct) = f'(x - ct)(-c).\end{aligned}$$

We might also do this problem as follows. Let $\xi(x, t) = x - ct$ and $\eta(x, t) = x + ct$. Then we have

$$u(x, t) = f(\xi(x, t)) + g(\eta(x, t)).$$

Now, a change in t causes changes in u through two processes, which we must sum over: (i) a change in ξ , which causes a change in f , and (ii) a change in η , which causes a change in g . Therefore, the chain rule applied to taking the derivative of u with respect to t gives us

$$\begin{aligned} \frac{\partial u}{\partial t} &= \frac{df}{d\xi} \frac{\partial \xi}{\partial t} + \frac{dg}{d\eta} \frac{\partial \eta}{\partial t} \\ &= f'(\xi)(-c) + g'(\eta)(+c) = -cf'(x - ct) + cg'(x + ct). \end{aligned}$$

Taking the other derivatives is a continuation of this process, and what has been presented in the first solution to this problem was essentially a short cut. Going through this process in a step-by-step fashion would yield the same result.

3. In equilibrium, since u will not change with respect to t , it must satisfy

$$\frac{d^2 u}{dx^2} = 0$$

by tossing out the partial t -derivative in the heat equation. Integrating this twice (note: this is now an ODE), we get

$$u(x) = Ax + B$$

for some constants A and B . (Interested students are recommended to also obtain this in a slightly different way – by considering the ODE as a constant-coefficient, second-order homogeneous ODE, and using the method of the characteristic equation.) Since we need $u(0) = T$, we have the condition $B = T$. Since we also need $u(L) = 0$, we get the condition $AL + B = 0$. Solving for A and B (our unknowns), we get $A = -T/L$ and $B = T$. Hence, the equilibrium temperature distribution is

$$u(x) = -\frac{T}{L}x + T.$$

4. The Q here represents heat sources, and the heat equation in this more general case is

$$c\rho \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left(K_0 \frac{\partial u}{\partial x} \right) + Q.$$

At equilibrium, there is no variation in t , and thus

$$0 = \frac{d}{dx} \left(K_0 \frac{du}{dx} \right) + Q.$$

Now substitute $Q = K_0 x^2$ as given in this problem, where K_0 be a constant. Thus we get the equation

$$\frac{d^2 u}{dx^2} = -x^2. \tag{1}$$

Integrating this once

$$\frac{du}{dx} = -\frac{1}{3}x^3 + A,$$

where A is an integration constant. Integrating again,

$$u = -\frac{x^4}{12} + Ax + B$$

where B is another constant. Alternatively, one can determine the general solution to (1) by using the ideas of constant-coefficient, second-order inhomogeneous equations. In any case, we now need to satisfy the condition $u(0) = T$. This gives $B = T$. Secondly, we need $u'(L) = 0$. This gives $A = L^3/3$. Thus the equilibrium temperature distribution in this case is

$$u(x) = -\frac{x^4}{12} + \frac{L^3 x}{3} + T.$$

5. Let $u_1(x, y)$ and $u_2(x, y)$ each be solutions to Laplace's equation as stated. Now consider

$$U(x, y) = C_1 u_1(x, y) + C_2 u_2(x, y)$$

for arbitrary constants C_1 and C_2 . Now,

$$\begin{aligned} \frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} &= C_1 \frac{\partial^2 u_1}{\partial x^2} + C_2 \frac{\partial^2 u_2}{\partial x^2} + \left(C_1 \frac{\partial^2 u_1}{\partial y^2} + C_2 \frac{\partial^2 u_2}{\partial y^2} \right) \\ &= C_1 \left(\frac{\partial^2 u_1}{\partial x^2} + \frac{\partial^2 u_1}{\partial y^2} \right) + C_2 \left(\frac{\partial^2 u_2}{\partial x^2} + \frac{\partial^2 u_2}{\partial y^2} \right) \\ &= 0. \end{aligned}$$

where each of the quantities in parentheses in the penultimate line above are zero since each of $u_1(x, t)$ and $u_2(x, t)$ satisfies Laplace's equation.

6. (a) Setting $u(r, t) = \phi(r) h(t)$ yields

$$\phi(r) \frac{dh}{dt} = \frac{k h}{r} \frac{d}{dr} \left(r \frac{d\phi}{dr} \right).$$

Dividing by $k \phi(r) h(t)$ yields

$$\frac{1}{k h} \frac{dh}{dt} = \frac{1}{r \phi(r)} \frac{d}{dr} \left(r \frac{d\phi}{dr} \right) = -\lambda$$

where λ is a separation constant, which must be independent of both r and t . This gives the two ODEs

$$\frac{dh}{dt} = -\lambda k h \quad \text{and} \quad \frac{1}{r} \frac{d}{dr} \left(r \frac{d\phi}{dr} \right) = -\lambda \phi.$$

This is of course alternatively written as

$$h'(t) = -\lambda k h(t) \quad \text{and} \quad \frac{1}{r} (r \phi'(r))' = -\lambda \phi(r).$$

Note that the r -equation is nasty – it is not constant-coefficient, since the independent variable r appears throughout. Thankfully, we are not asked to solve it!

(b) Set $u(x, y) = \phi(x) h(y)$, to give

$$h \frac{d^2 \phi}{dx^2} + \phi \frac{d^2 h}{dy^2} = 0.$$

Dividing by $\phi(x) h(y)$ yields

$$\frac{1}{\phi} \frac{d^2 \phi}{dx^2} = -\frac{1}{h} \frac{d^2 h}{dy^2} = -\lambda,$$

which separates out to the ODEs

$$\frac{d^2\phi}{dx^2} = -\lambda\phi \quad \text{and} \quad \frac{d^2h}{dy^2} = \lambda h.$$

Of course, you may have decided to write these in the form

$$\phi''(x) = -\lambda\phi(x) \quad \text{and} \quad h''(y) = \lambda h(y),$$

which would be equivalent.

(c) Set $u(x, t) = \phi(x) h(t)$, resulting in

$$\phi(x) h'(t) = k h(t) \phi'''(x).$$

Dividing by $k\phi(x)h(t)$,

$$\frac{1}{k h(t)} h'(t) = \frac{1}{\phi(x)} \phi'''(x) = \lambda.$$

This separates into the two ODEs

$$h'(t) = \lambda k h(t) \quad \text{and} \quad \phi'''(x) = \lambda \phi(x).$$

Note that you are not asked to *solve* any of the ODEs you get, which might be difficult – and in any case, you are not given any boundary/initial conditions, and so it cannot be done.