

Hamiltonian Noether Theorem

ϕ^s one parameter symmetry of H : $H \circ \phi^s = H$

If $\phi^s = \phi_K^s$ then K is an integral of H :

$$\dot{K} = \{K, H\} = 0, \quad \text{or} \quad K \circ \phi_H^t = K.$$

$$\frac{d}{ds} H \circ \phi_K^s = \{H, K\} = 0 \quad (\text{last week})$$

$$\Rightarrow \{K, H\} = 0$$

□

K is constant w.r.t. the flow of H .

H is constant w.r.t. the flow of K .

Reduction:

1) fix integral $E_K = \{x : K(x) = k\} = K^{-1}(k)$

2) identify points on the same orbit ϕ_K^s

1+2: reduction by 2 dimensions.

How to do that in practice?

Use invariants of ϕ_K^s as coordinates!

Recall week 1: \mathbb{R}^4 , $z_j = q_j + i p_j$, $j=1,2$

$$K = \frac{1}{2}|z_1|^2 + \frac{1}{2}|z_2|^2, \quad \phi_K^s(z_1, z_2) = (e^{is}z_1, e^{is}z_2)$$

where $H = f(s_1, s_2, s_3, s_4)$ s.t. $H \circ \phi_K^s = H$

$$s_1 = |z_1|^2, \quad s_2 = |z_2|^2, \quad s_3 = \operatorname{Re}(z_1 \bar{z}_2), \quad s_4 = \operatorname{Im}(z_1 \bar{z}_2)$$

Reduce to quotient $\Sigma_k = \{x : K(k) = k\} = S^3$
 $\Pi_k = \Sigma_k / S^1$ identify pts on the same ϕ_k^S orbit

Hopf map $\tilde{H} : \Sigma_k \rightarrow S^2$ $w_1 = q_1 q_2 + p_1 p_2$
 $w_2 = -q_1 p_2 + q_2 p_1$

$$w_1 = S_3, w_2 = S_4, w_3 = \frac{1}{2}(S_2 - S_1), w_4 = \frac{1}{2}(S_2 + S_1) = k$$

$$\{w_1, w_2\} = \frac{\partial w_1}{\partial q_1} \frac{\partial w_2}{\partial p_1} - \frac{\partial w_1}{\partial p_1} \frac{\partial w_2}{\partial q_1} + \frac{\partial w_1}{\partial q_2} \frac{\partial w_2}{\partial p_2} - \frac{\partial w_1}{\partial p_2} \frac{\partial w_2}{\partial q_2}$$

$$= q_2^2 - p_2(-p_2) + q_1(-q_1) - p_1^2$$

$$= q_2^2 + p_2^2 - q_1^2 - p_1^2 = (S_2 - S_1) = 2w_3$$

etc. ... $\{w_4, w_i\} = 0$

$$J_p = \begin{pmatrix} 0 & 2w_3 & -2w_2 \\ -2w_3 & 0 & 2w_1 \\ 2w_2 & -2w_1 & 0 \end{pmatrix}$$
 Poisson structure on \mathbb{R}^3 with coords w_1, w_2, w_3

J_p has Casimir $C = w_1^2 + w_2^2 + w_3^2$
 $J_p \nabla C = 0$, fix Casimir $\{C=c\} = S^2$

\tilde{H} is a Poisson map

reduced dynamics: $w = \begin{pmatrix} w_1 \\ w_2 \\ w_3 \end{pmatrix}$

$\dot{w} = J_p \nabla_w H$ Hamiltonian dynamical system on S^2

Remarks on Reduction

- invariants are usually related
- they may satisfy inequalities
 $s_i = |z_i|^2 \geq 0$
- reduced Hamiltonian is the original Hamiltonian written in invariants.

Example: Free rigid body (Euler top)

example with symmetry group $SO(3)$

Convention: X in body, x in space ($\in \mathbb{R}^3$)

$$x = R X, R \in SO(3), R R^T = I$$

$$\begin{aligned} \dot{x} &= \dot{R} X \\ &= \dot{R} R^T x = \hat{\omega} x = \omega \times x \end{aligned} \quad \begin{array}{l} \text{antisymmetric} \\ \text{matrix} \end{array} \quad X = \text{const.}$$

$$\omega = \begin{pmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{pmatrix}, \quad \hat{\omega} = \begin{pmatrix} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{pmatrix}$$

$$\omega = R \Omega$$

$$\begin{aligned} \text{claim: } \hat{\Omega} &= R^T \hat{\omega} \\ \text{Identity: } \hat{R} a &= R \hat{a} R^T \end{aligned} \quad \left. \begin{array}{l} \hat{\Omega} = R^T \hat{\omega} \\ = R^T \hat{\omega} R \\ = R^T \dot{R} R^T R \\ = R^T \dot{R} \end{array} \right\}$$

$$\begin{aligned} \hat{R} a b &= (R a) \times b \\ &= R (a \times R^T b) \\ &= R \hat{a} R^T b \quad \checkmark \end{aligned}$$

Configuration space $SO(3) \ni R$

$$\phi^g(R, \dot{R}) = (gR, g\dot{R}), \quad g \in SO(3)$$

observe: $\hat{S} \cdot \phi^g = \hat{S}$ invariant

$$(gR)^T g \dot{R} = R^T g^T g \dot{R} = R^T \dot{R} \quad \checkmark$$

Tensor of inertia $\Theta, \quad \Theta^T = \Theta$

$L = \Theta \Omega$ angular momentum, invariant

equations of motion: $L = \begin{pmatrix} L_1 \\ L_2 \\ L_3 \end{pmatrix}$ $J_E = L$ ϕ^g

$$\dot{L} = L \times \Omega = L \times \Theta^{-1} \dot{L} = J_E \nabla_L H, \quad H = \frac{1}{2} (L \Theta^{-1} L) \quad (\text{kinetic energy})$$

$$\text{finding } R: \quad \hat{S} = R^T \dot{R} \Rightarrow \dot{R} = R \hat{\Omega}$$

Θ diagonal $= \text{diag}(I_1, I_2, I_3)$
 $I_1 > I_2 > I_3$

$$\begin{pmatrix} \dot{L}_1 \\ \dot{L}_2 \\ \dot{L}_3 \end{pmatrix} = \begin{pmatrix} L_1 \\ L_2 \\ L_3 \end{pmatrix} \times \begin{pmatrix} L_1/I_1 \\ L_2/I_2 \\ L_3/I_3 \end{pmatrix} = \begin{pmatrix} L_2 L_3 (\frac{1}{I_2} - \frac{1}{I_3}) \\ L_1 L_3 (\frac{1}{I_3} - \frac{1}{I_1}) \\ L_1 L_2 (\frac{1}{I_1} - \frac{1}{I_2}) \end{pmatrix} = \begin{pmatrix} \alpha L_2 L_3 \\ -\beta L_1 L_3 \\ \gamma L_1 L_2 \end{pmatrix}$$

$$\alpha, \beta, \gamma \geq 0$$

Equilibria:

$$L_1 = L_2 = 0: \quad L = \mu \hat{e}_3$$

$$L_1 = L_3 = 0: \quad L = \mu \hat{e}_2$$

$$L_2 = L_3 = 0: \quad L = \mu \hat{e}_1$$

lin. stability

$$DJ_E^{\nabla H} = \begin{pmatrix} 0 & \alpha L_3 & \alpha L_2 \\ -\beta L_3 & 0 & -\beta L_1 \\ \gamma L_2 & \gamma L_1 & 0 \end{pmatrix}$$

$$L = \mu \hat{i} \quad D^2 \nabla H(\mu \hat{i}) = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -\mu\beta \\ 0 & \mu\gamma & 0 \end{pmatrix}$$

char. poly. $\lambda(\lambda^2 + \mu^2\beta\gamma) \Rightarrow$ lin. stable

$L = \mu \hat{j} \quad \dots \quad \lambda(\lambda^2 - \mu^2\alpha\gamma) \Rightarrow$ unstable

$L = \mu \hat{k} \quad \lambda(\lambda^2 + \mu^2\alpha\beta) \Rightarrow$ lin. stable

Casimir $L_1^2 + L_2^2 + L_3^2 = \mu^2$ sphere

Holland $\frac{L_1^2}{I_1} + \frac{L_2^2}{I_2} + \frac{L_3^2}{I_3} = 2h$ ellipsoid

