

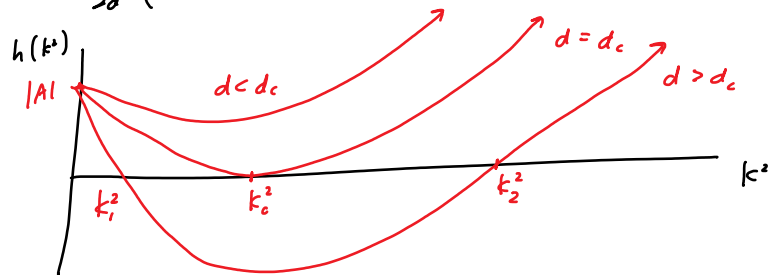
When  $d > d_c$ , the range of unstable wave numbers  $k$  is constrained by

$$k_1^2 < k^2 < k_2^2,$$

where  $k_1^2$  and  $k_2^2$  are the zeros of  $h(k^2) = 0$ , i.e.,

$$k_1^2 = \frac{\gamma}{2d} \left( df_u + g_v - \sqrt{(df_u + g_v)^2 - 4d|A|} \right),$$

$$k_2^2 = \frac{\gamma}{2d} \left( \text{''} + \text{''} \right).$$



Then, the solution

$$\vec{w} = \sum_{\mathbf{k}} c_{\mathbf{k}} e^{\lambda(\mathbf{k}^2)t} \vec{W}_{\mathbf{k}}(\vec{r})$$

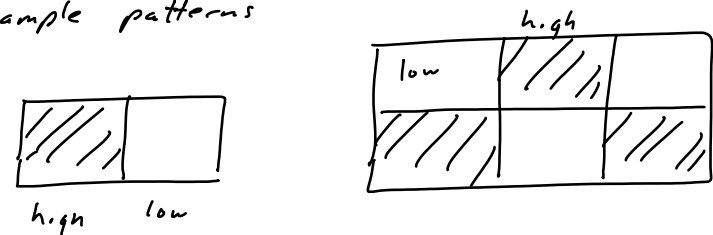
approaches

$$\sum_{\mathbf{k}} c_{\mathbf{k}} e^{\lambda(\mathbf{k}^2)t} \vec{W}_{\mathbf{k}}(\vec{r})$$

for large  $t$ .

Note that the system is nonlinear, so a key assumption is that these linearly unstable modes will eventually be bounded by nonlinear terms, resulting in a spatially inhomogeneous steady state.

Example patterns



In summary, the conditions for spatial pattern formation for

$$u_t = \gamma f(u, v) + \nabla^2 u,$$

$$v_t = \gamma g(u, v) + d \nabla^2 v$$

with no-flux boundary conditions are

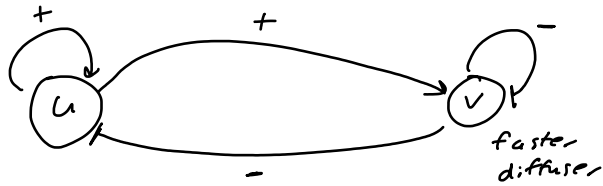
- ①  $f_u + g_v < 0$
- ②  $f_u g_v - f_v g_u > 0$
- ③  $df_u + g_v > 0$
- ④  $(df_u + g_v)^2 - 4d(f_u g_v - f_v g_u) > 0$

From ① and ③,  $f_u$  and  $g_v$  must have opposite signs. Suppose  $f_u > 0$  and  $g_v < 0$ , then ③  $\Rightarrow d > 1$ .

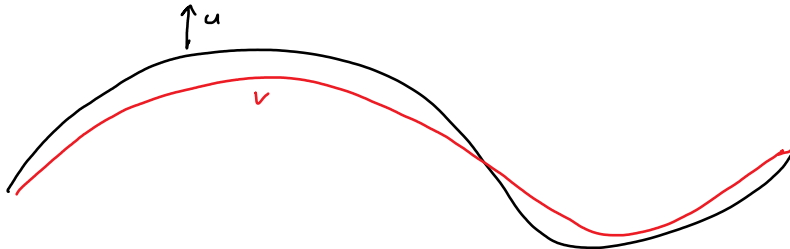
Since  $f_u g_v < 0$ , ②  $f_v g_u < f_u g_v < 0$ , so  $f_v$  and  $g_u$  also have opposite signs.

We have 2 cases:

$$i) \begin{bmatrix} f_u & f_v \\ g_u & g_v \end{bmatrix} = \begin{bmatrix} + & - \\ + & - \end{bmatrix}$$

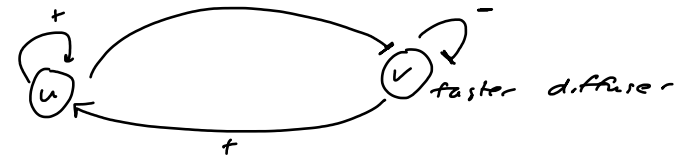


$u$  is the activator of  $v$  (and itself), and  $v$  is the inhibitor of  $u$  (and itself).



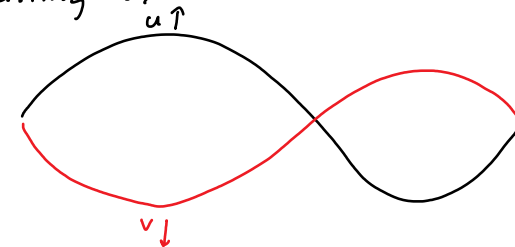
They are in phase.

$$ii) \begin{bmatrix} f_u & f_v \\ g_u & g_v \end{bmatrix} = \begin{bmatrix} + & + \\ - & - \end{bmatrix}$$



$u$  is the inhibitor of  $v$ , and  $v$  is the activator of  $u$ .

Resulting dynamics look like



They are out of phase

Specific example

Consider the 1-D system

$$u_t = \delta f(u, v) + u_{xx},$$

$$v_t = \delta g(u, v) + d v_{xx},$$

where  $f(u, v) = a - u + u^2v,$   
 $g(u, v) = b - u^2v,$

for some positive constants  $\delta, a, b, d.$

The positive steady state is

$$u_0 = a + b,$$

$$v_0 = \frac{b}{(a+b)^2}.$$

At steady state,

$$f_u = \frac{b-a}{a+b}$$

$$f_v = (a+b)^2 > 0$$

$$g_u = -\frac{2b}{a+b} < 0$$

$$g_v = -(a+b)^2 < 0.$$

Since  $f_u$  and  $g_v$  must have opposite signs, we must have  $f_u > 0 \Rightarrow b > a.$

From the 4 "Turing conditions", we have

$$f_u + g_v < 0 \Rightarrow b - a < (a+b)^3$$

and

$$f_u g_v - f_v g_u > 0 \Rightarrow (a+b)^2 > 0 \quad \text{automatically satisfied}$$

and

$$d f_u + g_v > 0 \Rightarrow d(b-a) > (a+b)^3$$

and

$$(d f_u + g_v)^2 - 4d(f_u g_v - f_v g_u) > 0$$

$$\Rightarrow (d(b-a) - (a+b)^3)^2 > 4d(a+b)^4$$

These conditions define a domain in  $(a, b, d)$ -parameter space, called the Turing space, in which patterns from diffusion-driven instability can occur.

Consider the domain  $x \in (0, L)$  and consider the eigenvalue problem

$$W_{xx} + k^2 W = 0$$

with BCs  $\frac{\partial W(0, t)}{\partial x} = \frac{\partial W(L, t)}{\partial x} = 0.$

The solutions are

$$W_n(x) = A_n \cos\left(\frac{n\pi x}{L}\right)$$

where  $A_n$  is constant.

So, the wave numbers are  $k = \frac{n\pi}{L}$ .

Following the same reasoning as for the general case, we get

$$k_1^2 < k^2 < k_2^2$$

where

$$k_1^2 = \gamma \left( \frac{d(b-a) - (a+b)^3 - \sqrt{(d(a-b) - (a+b)^3)^2 - 4d(a+b)^4}}{2d(a+b)} \right)$$

$$k_2^2 = \gamma \left( \frac{\text{"} + \text{"}}{\text{"}} \right)$$

In terms of wavelength,

$$\omega = \frac{2\pi}{\lambda}$$

and  $\omega_1^2 < \omega^2 < \omega_2^2$  where  $\omega_i = \frac{2\pi}{\lambda_i}$

Note that both sides of the inequalities for  $k^2$  are proportional to  $\gamma$ . Also, the allowed wave numbers  $k = \frac{n\pi}{L}$  are discrete.

So, for small  $\gamma$ , there are no unstable modes.

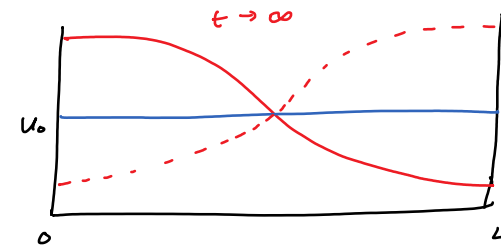
Suppose  $\gamma$  is such that the only unstable wave number corresponds to  $n=1$ .

Then,

$$\vec{w}(x,t) \sim \left( \exp\left(\lambda \left(\frac{\pi^2}{L^2}\right) t\right) \cos \frac{\pi x}{L} \right)$$

as  $t \rightarrow \infty$ .

So, one would expect a spatial pattern like



Recall that unstable modes satisfy

$$k_1^2 < k^2 < k_2^2$$

where  $k_1^2 = \gamma(\text{"} - \text{"}) = \gamma P$   
 $k_2^2 = \gamma(\text{"} + \text{"}) = \gamma Q$

and  $k = \frac{n\pi}{L}$ .

$$\text{So, } \gamma P \left(\frac{L}{\pi}\right)^2 < n^2 < \gamma Q \left(\frac{L}{\pi}\right)^2$$

and suppose  $n=1$  is the only  $n \in \mathbb{N}$  that satisfies this inequality.

Suppose we double the spatial domain size.  
Then  $\gamma$  will change to  $4\gamma$ .

Then  $n=2$  satisfies

$$(4\gamma) \left(\frac{L}{\pi}\right)^2 \rho < n^2 < (4\gamma) \left(\frac{L}{\pi}\right)^2 Q,$$

so the  $n=2$  mode can arise

