Tuesday, September 19, 2017 9:59 AM

Discussion Question 4

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 $\omega = \alpha - \omega_{o}$ 

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1.

$$\begin{aligned} \frac{\partial u}{\partial t} &= \mathcal{D} \frac{\partial u}{\partial x^{2}} + f(u) & \text{no-flux } \theta C_{3} \\ u_{t} &= f(u) & u_{0} & 2 \cdot u_{t} = \mathcal{D} \frac{\partial^{2} u}{\partial a^{2}} + f(u) \\ \vdots & u_{t-u_{0}} & u_{t} = \mathcal{D} \frac{\partial^{2} u}{\partial a^{2}} + f_{u}(u_{0})w - (A) \\ & u_{t} &= f_{u}(u_{0})w \\ & f_{u}(u_{0}) < o & u(x_{1}t) = \sum_{k} c_{k} \frac{e^{A_{k}t} W_{k}(x_{1})}{\mathcal{D} \frac{\partial^{2} W_{k}}{\partial x^{2}} + k^{2}W_{k} = 0} - (B) \\ & \frac{\partial}{\partial t} \left( e^{A_{k}t} W_{k}(x_{1}) = f_{u}(u_{0}) e^{A_{k}t} W_{k}(x_{1}) + \mathcal{D} \frac{\partial^{2}}{\partial x^{2}} (e^{A_{k}t} W_{k}(x_{1})) \\ & \lambda_{k} e^{A_{k}t} W_{k}(x_{1}) = f_{u}(u_{0}) e^{A_{k}t} W_{k}(x_{1}) + \mathcal{D} \frac{\partial^{2}}{\partial x^{2}} (e^{A_{k}t} W_{k}(x_{1})) \\ & \frac{\partial u}{\partial x^{2}} = -k^{2}W_{k}(x_{1}) \\ & \lambda_{k} = f_{u}(u_{0}) + -k^{2} \\ & Re(\lambda_{k}) > 0 & -) & f_{u}(u_{0}) - 4k^{2} = 0 \\ & \Rightarrow & k^{2} < f_{u}(u_{0}) < 0 & \frac{1}{2} \\ & \frac{\partial u^{2}}{\partial x^{2}} \end{aligned}$$

For convenience, let  

$$\hat{g}(x,t) = g(X,t) - g$$
,  
and drop the  $\wedge$  notation.  
Also, yeast moves very slowly compared to the rate  
of glucose diffusion, so let us assume that yeasy  
is non-mobile, i.e.  $D_e = O$ .  
So, we remark the equations as  
 $\frac{\partial n}{\partial t} = Eng$   
 $\frac{\partial g}{\partial t} = D \frac{\partial^2 g}{\partial x^2} - cEng$   
Let  $N(2) = n(x,t)$  and  $\hat{b}(2) = g(x,t)$  where  $Z = x - vt$ .  
Then,  $-v \frac{dN}{dz} = ENG$   
 $-v \frac{dG}{dz} = D \frac{d^2 G}{dz^2} - cENG$   
We could turn this into a 3-D system, or we  
could take  $C \cdot (1 + (2))$   
 $\rightarrow -vc \frac{dN}{dz} - v \frac{dG}{dz} = D \frac{d^2 G}{dz^2}$ 

Integrating, we get  

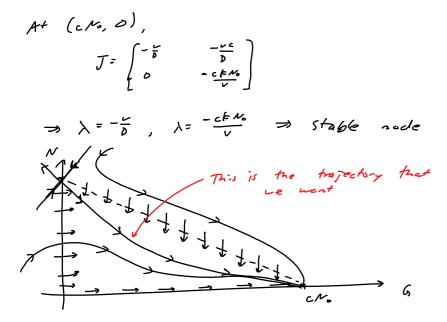
$$\int_{-\infty}^{2} \left( -vc \frac{du}{dz} - v \frac{d\theta}{dz} \right) dz = \int_{-\infty}^{2} \left( \frac{d^{2}\theta}{dz^{2}} + dz \right) \\
-vc N - v \left( \frac{d}{dz} \right) - v \left( \frac{d}{dz} \right) + vc N \left( - \alpha \right) + v \left( \frac{d}{dz} \right) - \frac{d}{dz^{2}} - \frac{d}{dz^{2}} \right) \\
-vc N(z) - v G(z) + vc N \left( - \alpha \right) + v G(-\alpha 0) = \frac{d}{dz^{2}} - \frac{d}{dz^{2}}$$

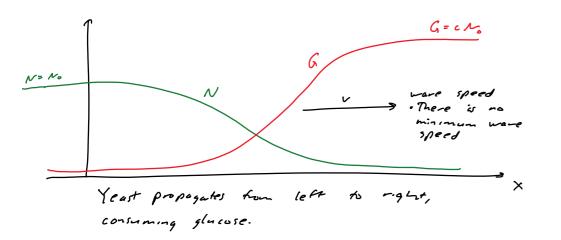
$$J_{acobian} \qquad J = \begin{bmatrix} -\frac{v}{0} & -\frac{vc}{0} \\ -\frac{k}{v} & -\frac{kG}{v} \end{bmatrix}.$$

$$A + (0, N_0),$$

$$det (J - \lambda I) = \lambda^2 + \frac{\omega}{D}\lambda - \frac{kcN_0}{D} = 0$$

$$\Rightarrow \lambda = \frac{-\frac{\omega}{D} + \sqrt{\frac{\omega^2}{D^2} + \frac{4kcN_0}{D}}}{2} \Rightarrow saddle$$





Travelling Pulses A travelling wave solution that starts and ends at the some steady state of the governing equations. Fitz Hugh - Naguno Equations  $\mathcal{E}\frac{\partial v}{\partial t} = \mathcal{E}^2 \frac{\partial^2 v}{\partial x^2} + \mathcal{F}(v,w)$  $\frac{\partial w}{\partial t} = g(u, w)$ where E>O is assumed to be very small. Apply travelling wore courdinates Z=X-ct, where CZO. we get  $E^2 V_{22} + CE V_2 + f(U, w) = 0$ CW2 + g(4, w) = 0. Let us examine the simplified, preceivise linear f(r,w) = H(r-w) - v - wcase g(u, w)= U where  $H(x) = \begin{cases} 1, & x \ge 0 \\ 0, & x \ge 0 \end{cases}$ In this case, we can construct exact solutions.

The steady state where f(yw) = g(y,w)=0 is (0,0).